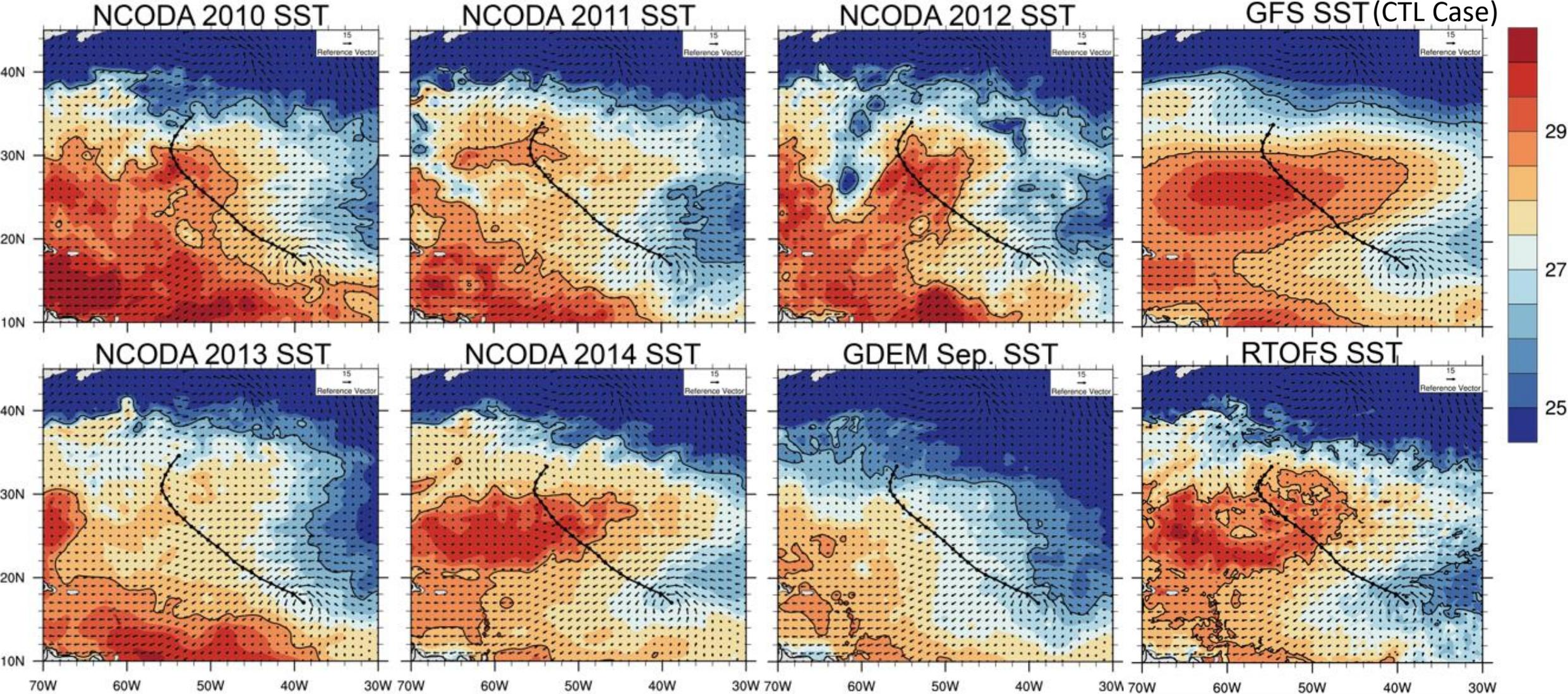


# An analysis of Hurricane Edouard SST sensitivity runs by HWRF in a neutral-shear environment

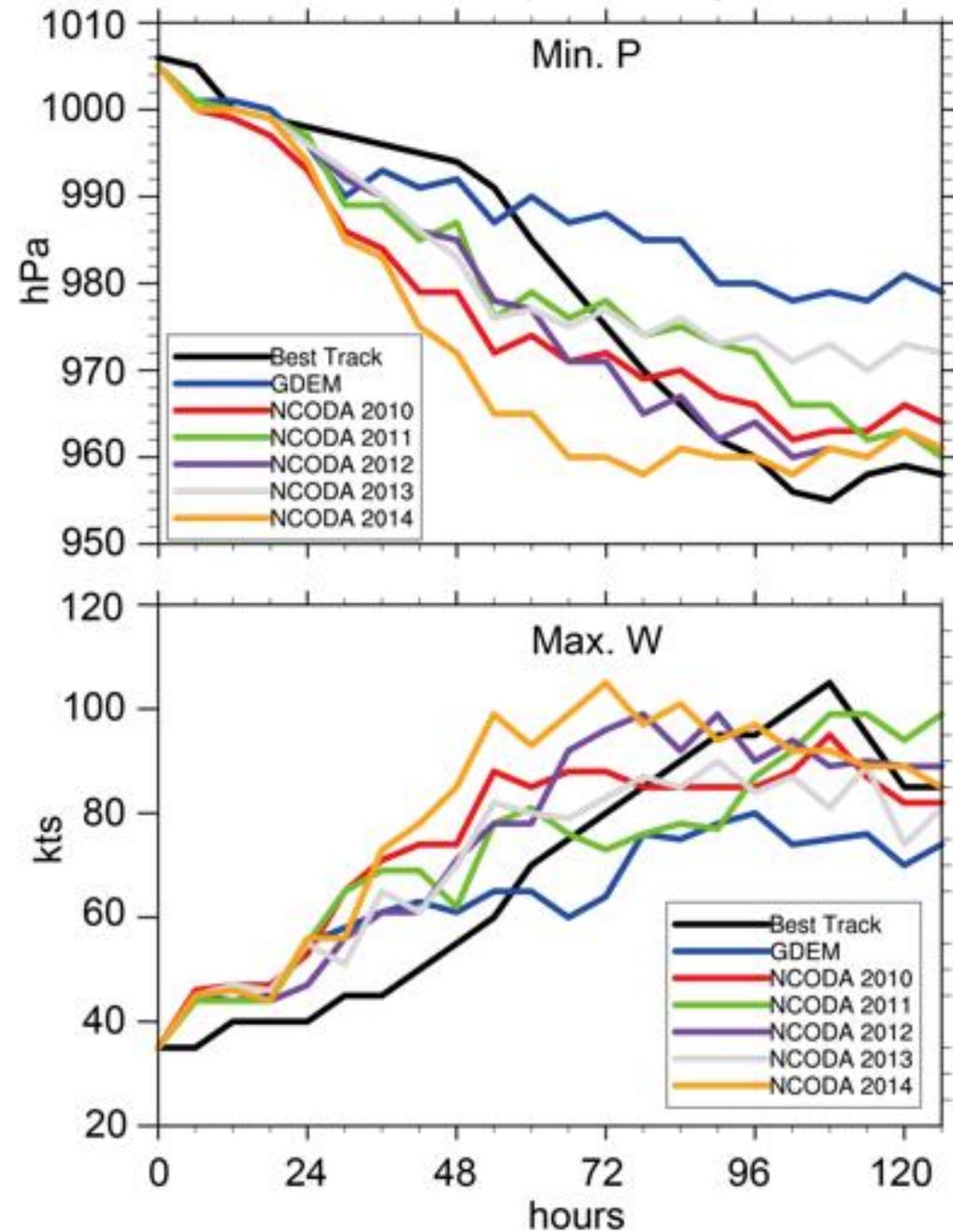
Pat Fitzpatrick<sup>3</sup>, Yee Lau<sup>3</sup>, Jili Dong<sup>1,2</sup>, Lin Zhu<sup>1,2</sup>, Hyun-Sook Kim<sup>1,2</sup>, Avichal Mehra<sup>1</sup>  
<sup>1</sup>NOAA/NWS/NCEP/EMC, <sup>2</sup>IMSG, Inc., <sup>3</sup>Mississippi State University

- Review of original sensitivity runs
- Shear calculation issues; modified Kurihara filter
- Time series analysis for shear < 14 m/s
  - Goal is to assess if we can obtain a basic understanding of the intensity changes and steady-state conditions in sensitivity run performed with HWRF for Hurricane Edouard during favorable environmental conditions (i.e., low-to-neutral shear, moist background environment).
- Relationship analysis to  $V_{\max}$  and 24-hr intensity change
  - a. Also looked at 6-h and 12-h intensity change, results similar, lower correlations.
- Maximum Potential Intensity applications?

# SST fields for HWRF simulations in Edouard (2014) environment



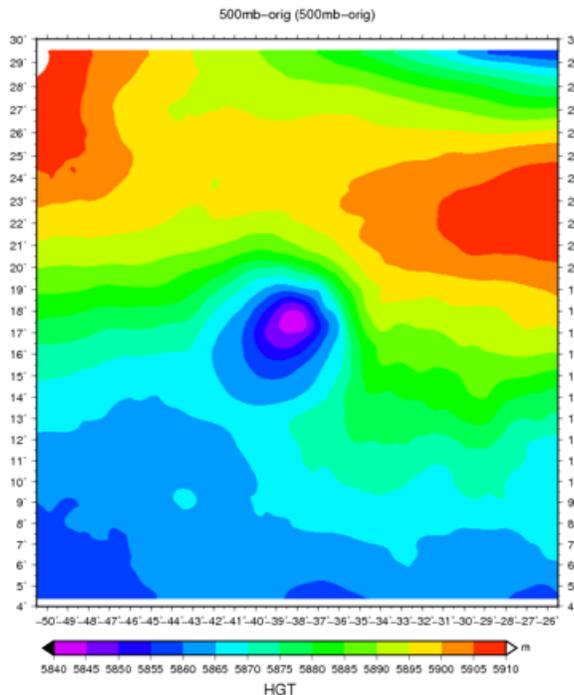
# Intensity change



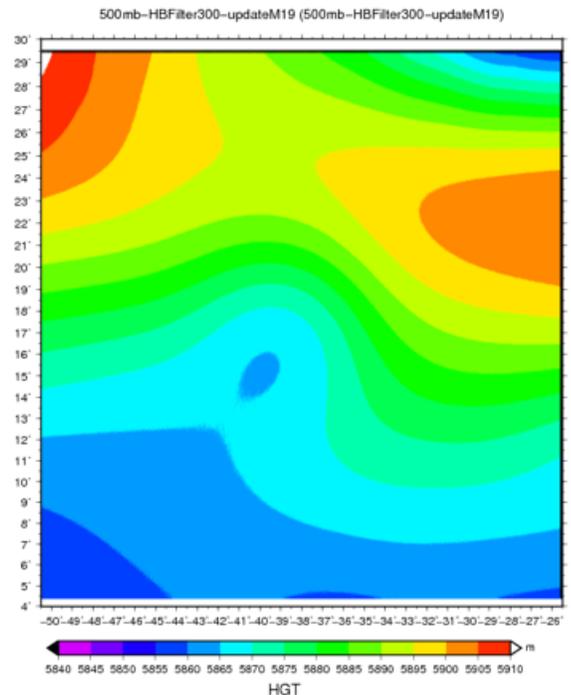
# Shear calculation required filtering vortex

- Kurihara filter used to remove vortex (see below)
- We expanded the original Kurihara filter with the following weights:  
2,3,4,2,5,6,7,2,8,9,2,10,11,2,12,13,2,3,2
- A periodic 2 is needed to stop numerical instabilities from the filter. More passes were required than the original scheme as well.
- Response function derivation and analysis available upon request

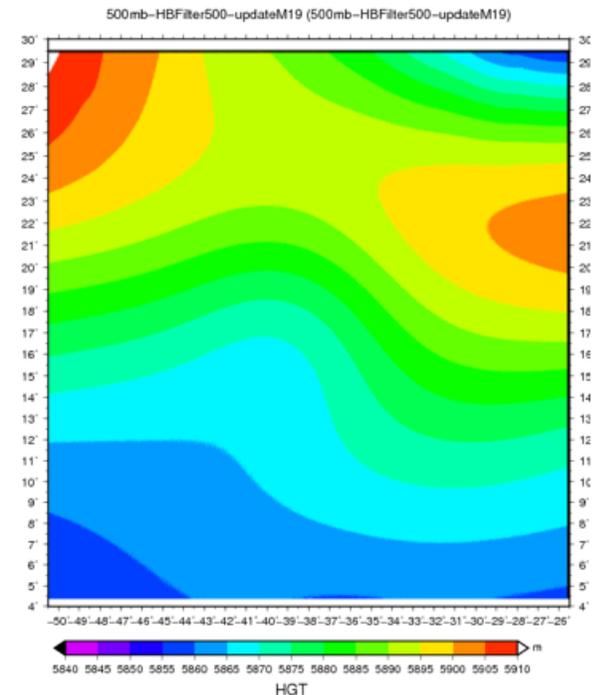
**Original**



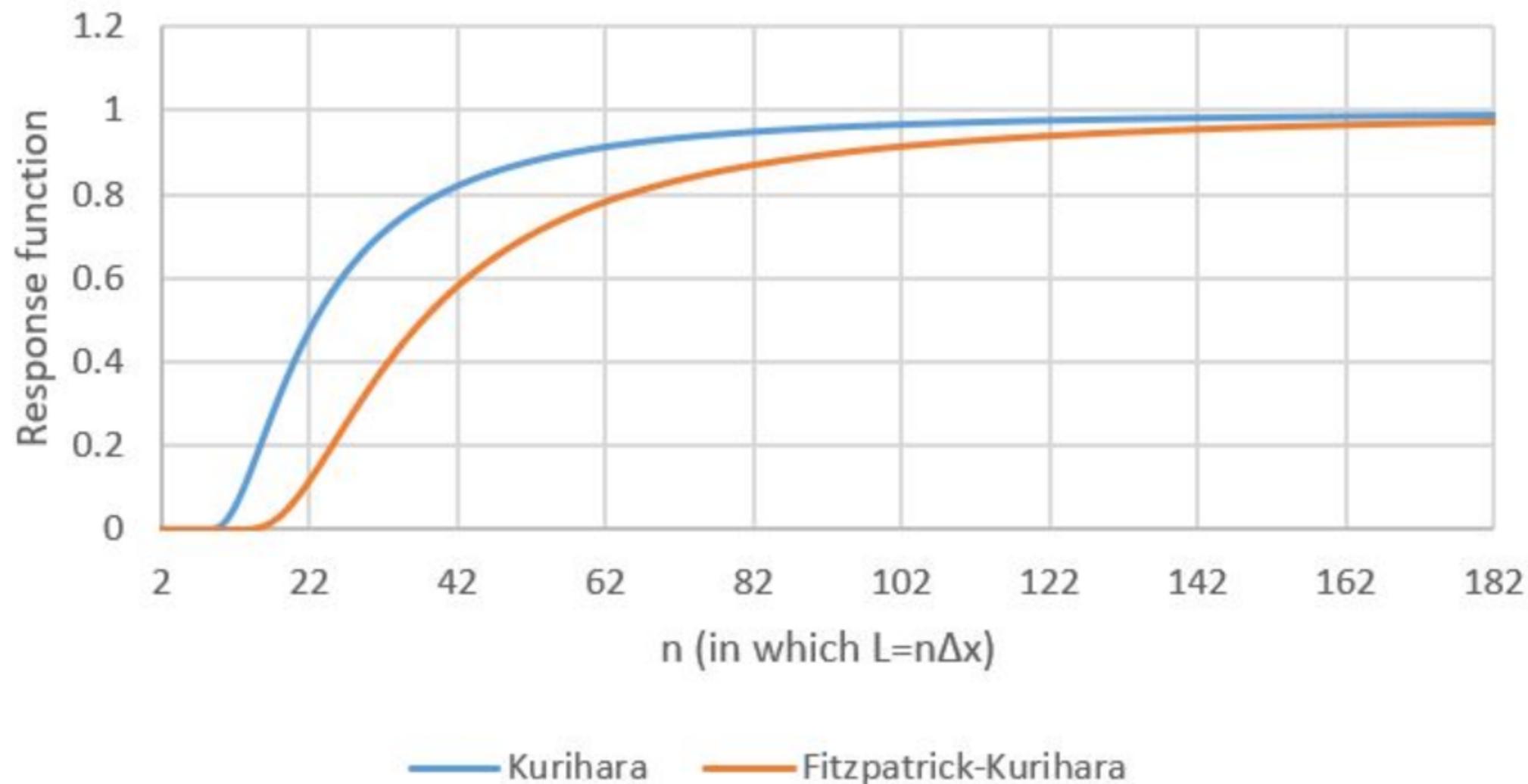
**300 passes  
Enhanced Kurihara filter**



**500 passes  
Enhanced Kurihara filter**



## Attenuation per pass

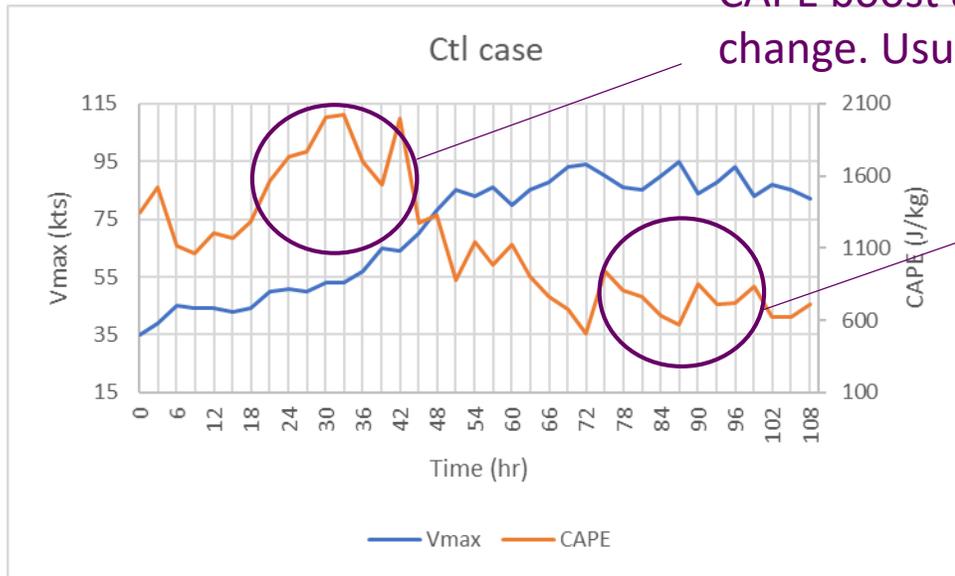


## Time series analysis

Looked at fluxes, CAPE, dewpoint, and RH at Rmax averaged every 30 compass degrees, resulting in 12-point average.

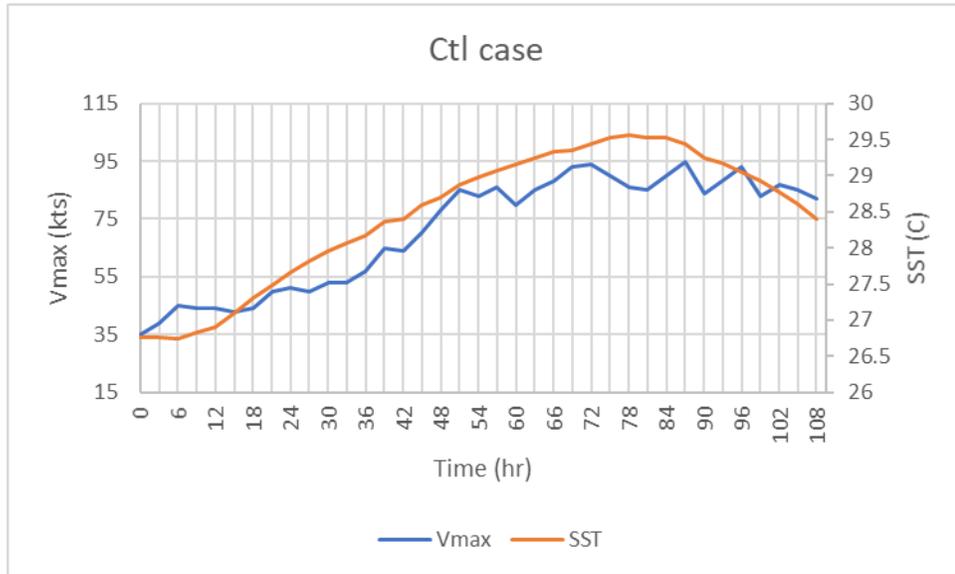
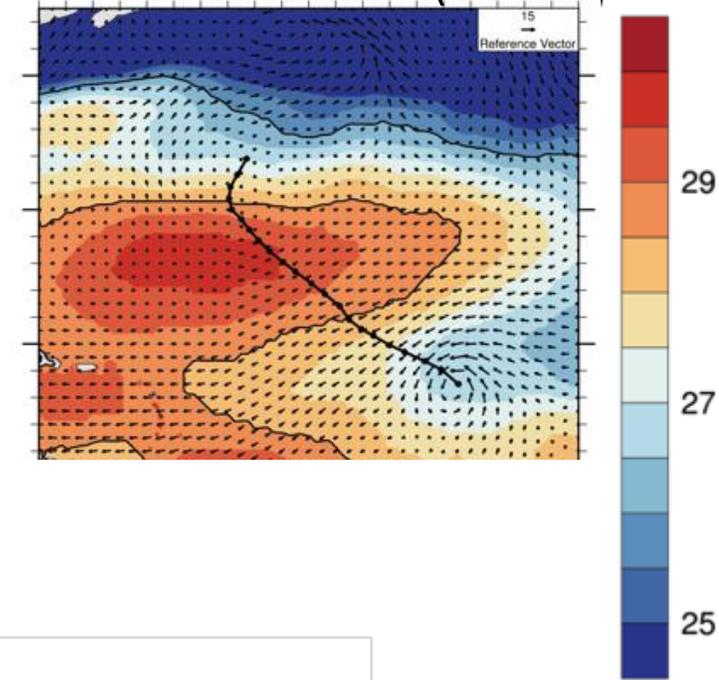
Looked at filtered wind shear and PW at 100, 200, 300, 400, and 500-km radii in 90 compass degrees, resulting in 20-point average

CAPE boost associated with faster intensity change. Usually when SST 27.5-28.5 C

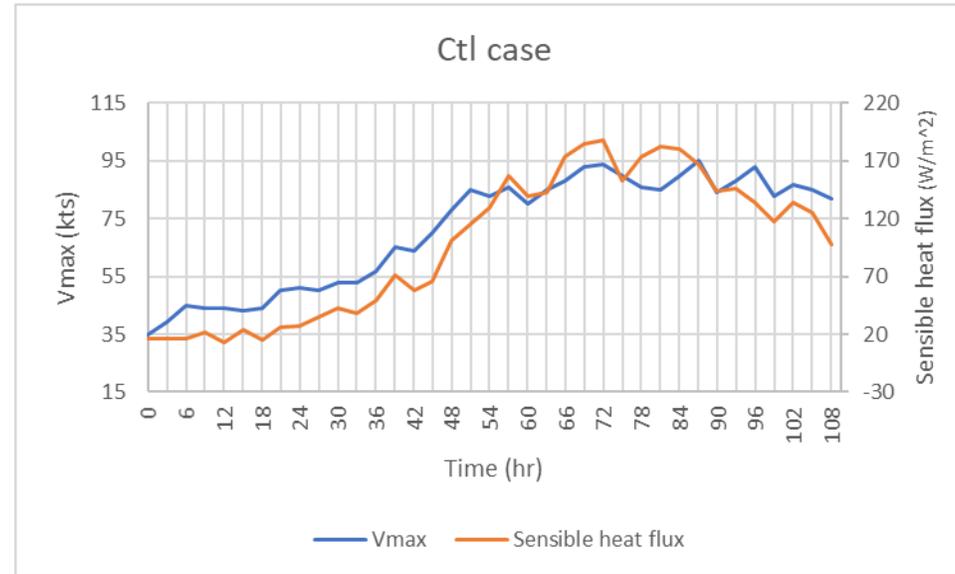


Steady state associated with relatively small CAPE

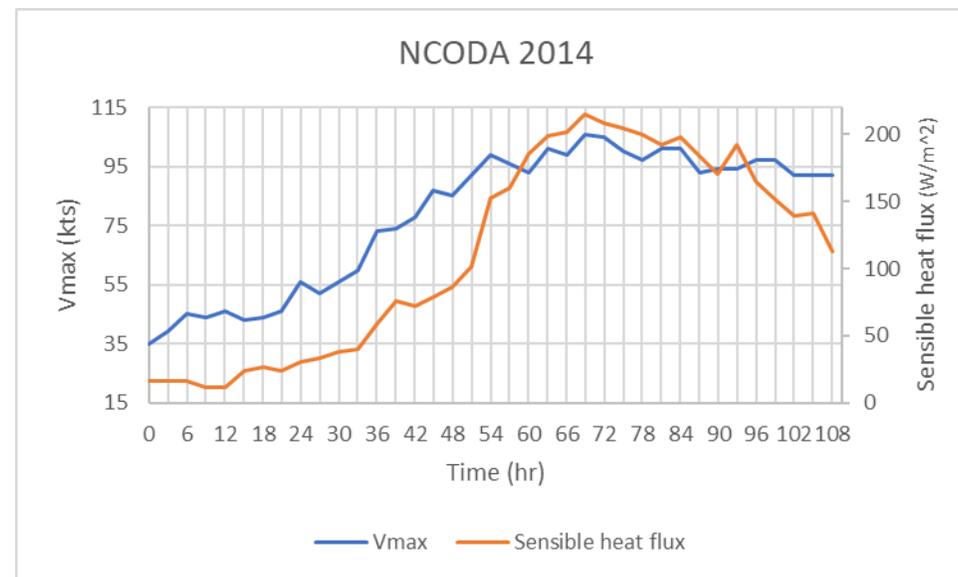
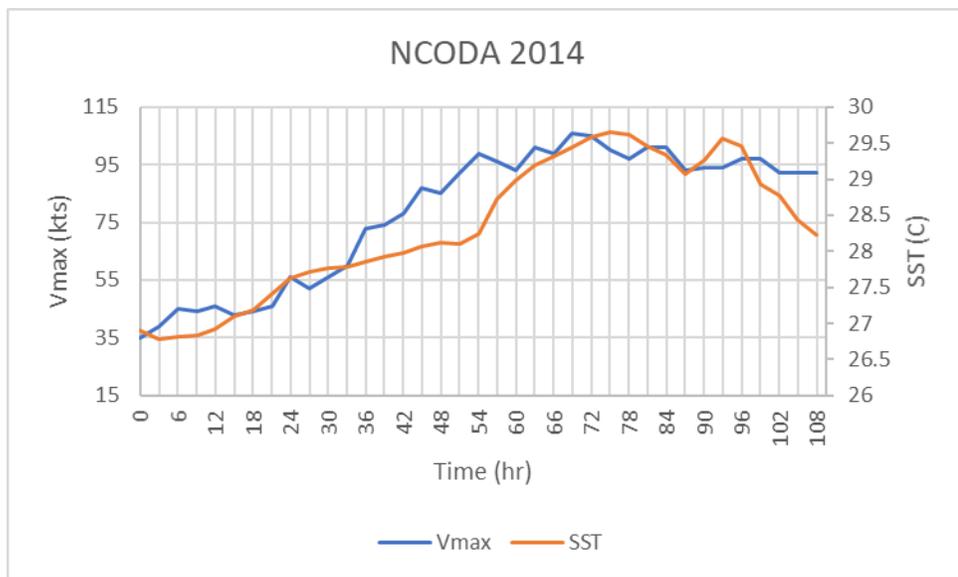
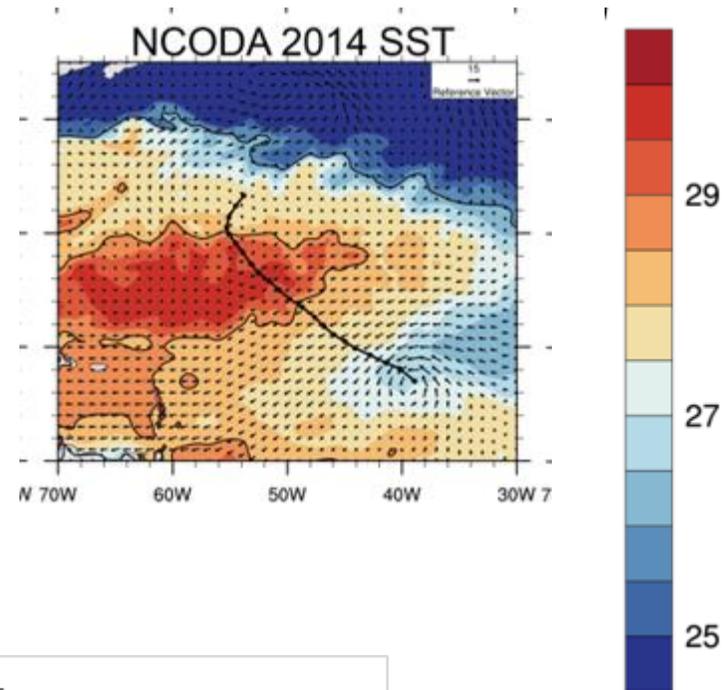
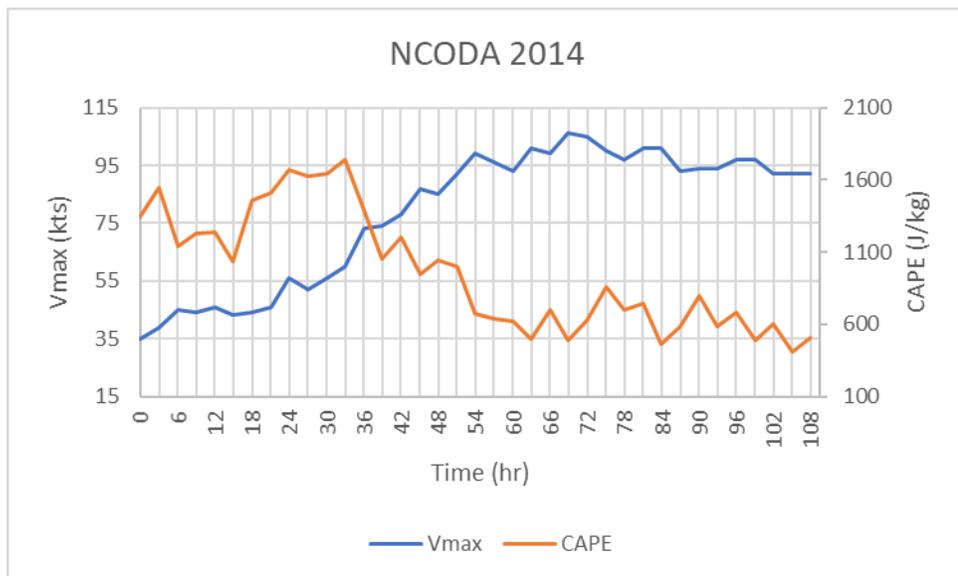
GFS SST (CTL Case)

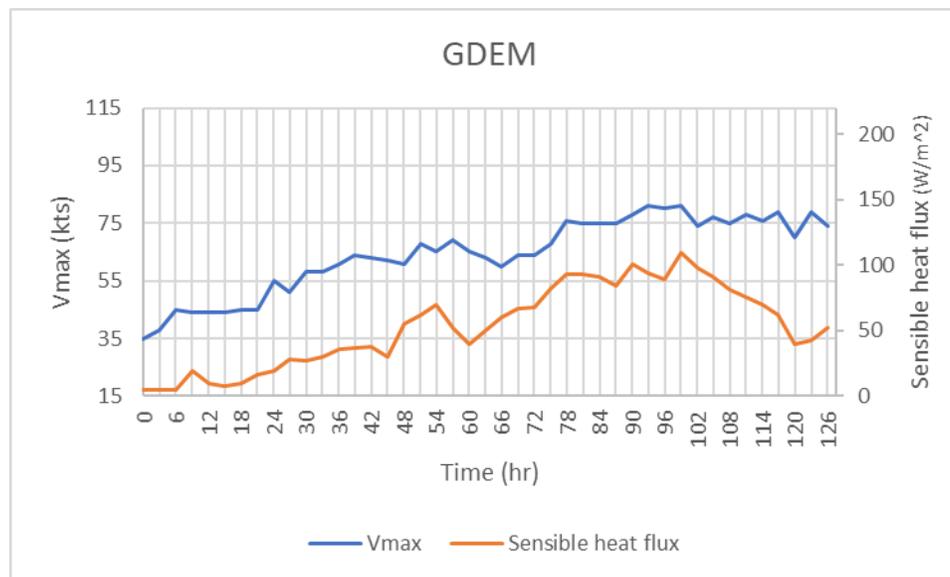
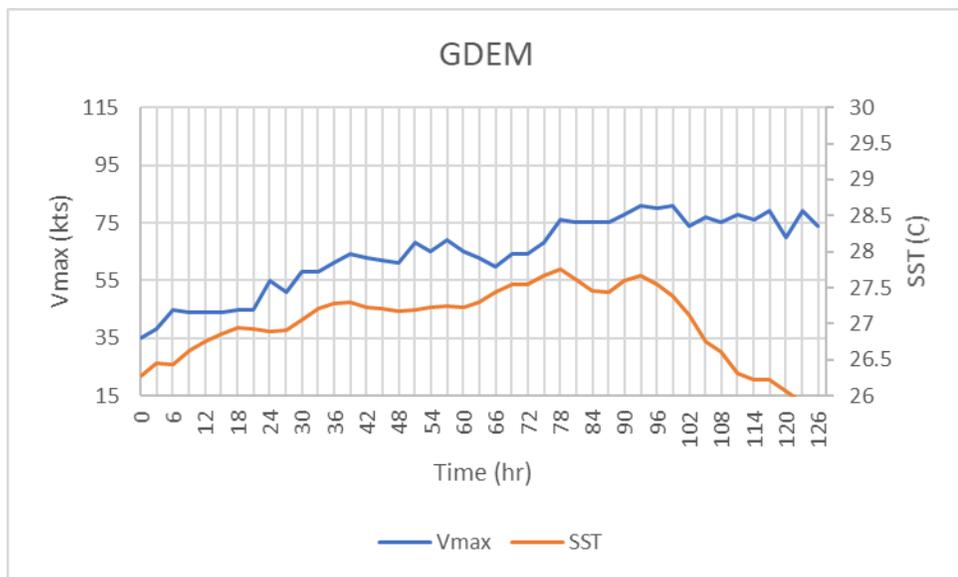
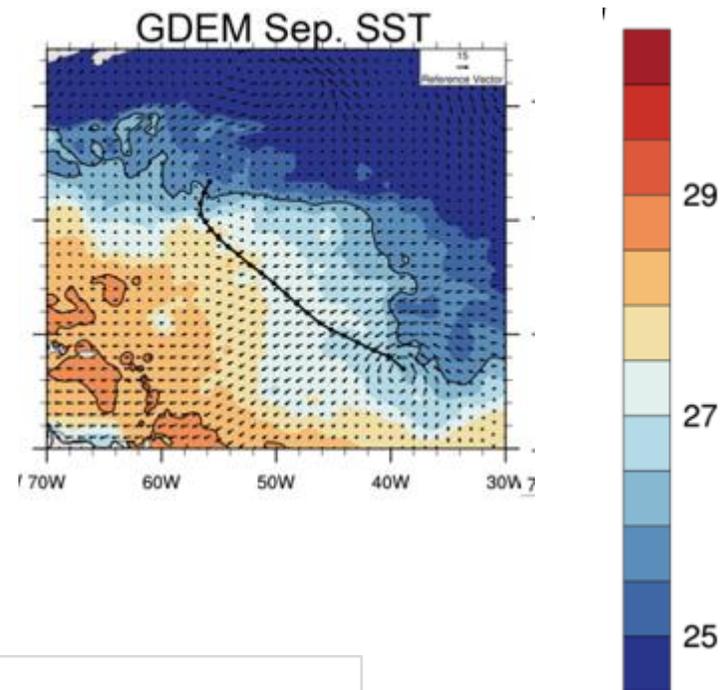
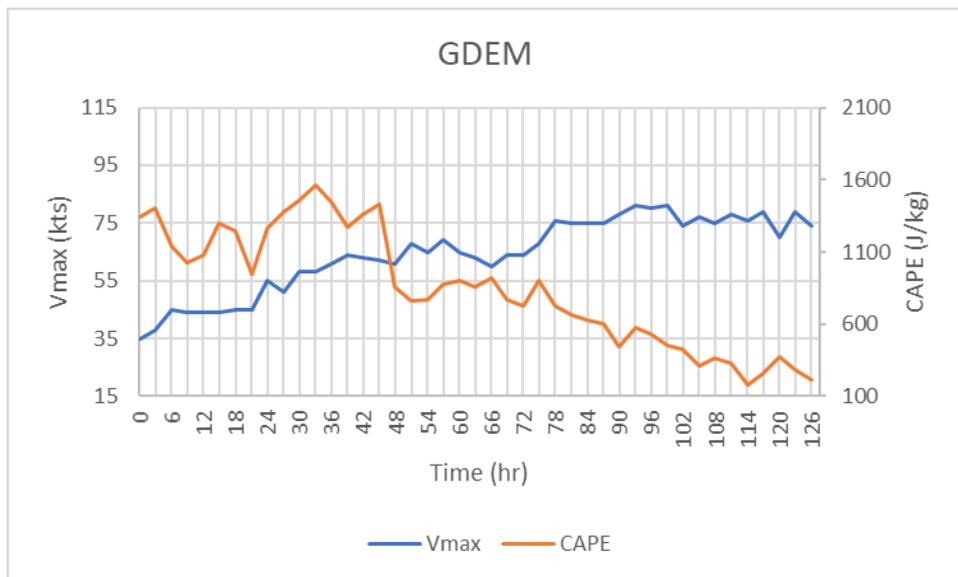


Linear relationship of  $V_{max}$  to SST obvious  
 $r=0.77$  to instantaneous SST,  $n=313$  ( $r^2=58.8\%$ )  
 $r=0.78$  for 24-hr avged SST,  $n=280$  ( $r^2=60.9\%$ )

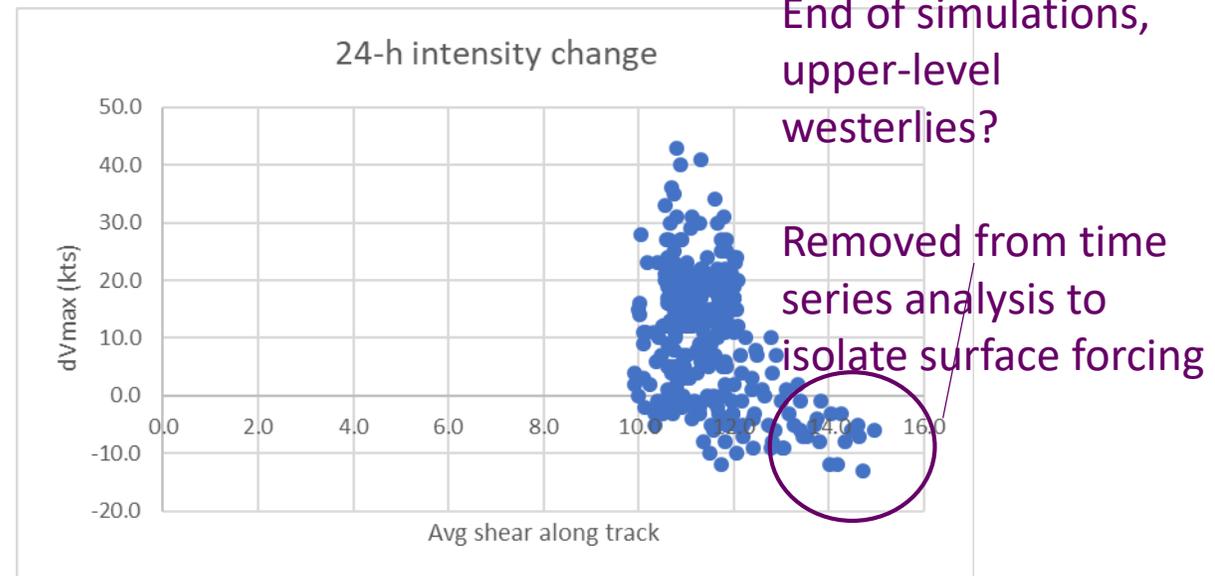
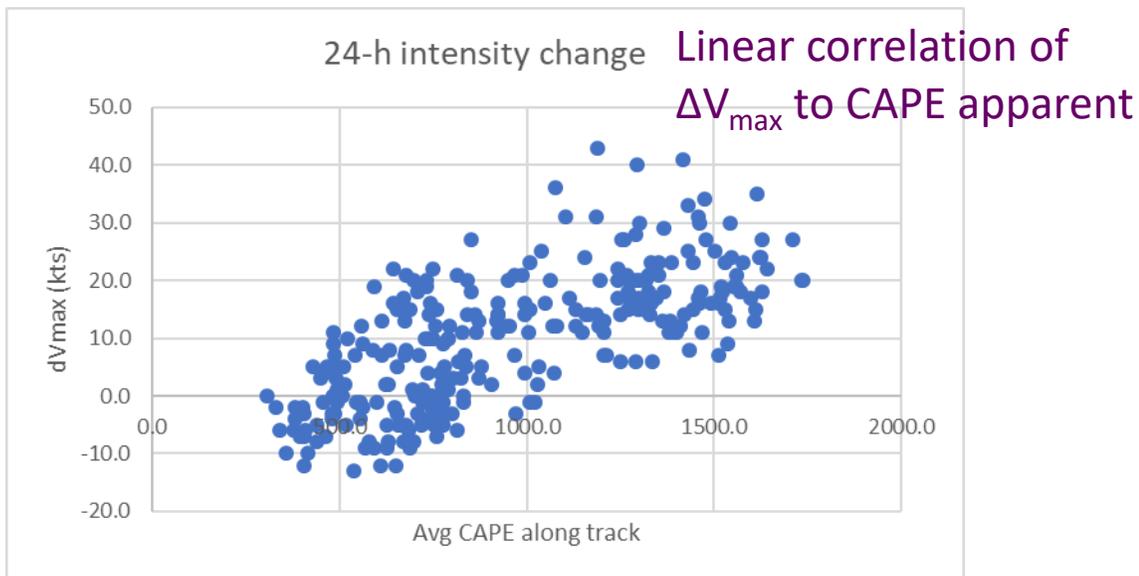
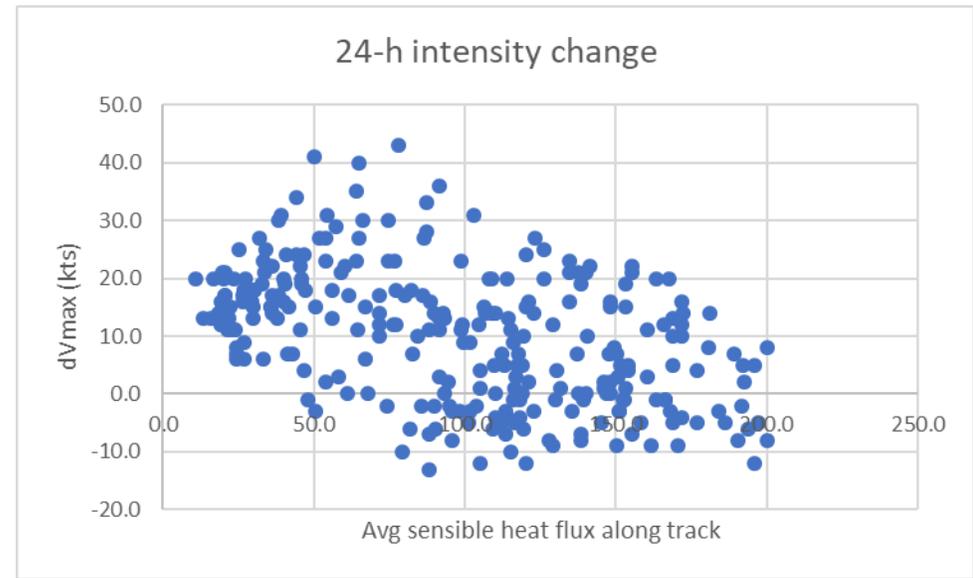
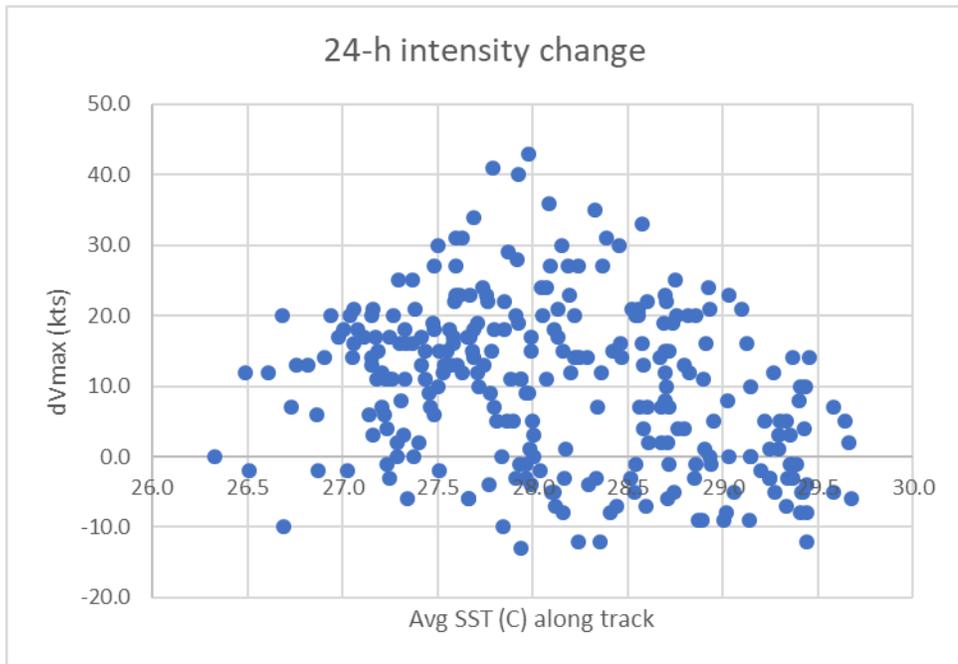


Linear relationship of  $V_{max}$  to heat flux higher than SST  
 $r=0.89$  to instantaneous heat flux,  $n=313$  ( $r^2=79.6\%$ )  
 $r=0.85$  for 24-hr avged heat flux,  $n=280$  ( $r^2=71.6\%$ )



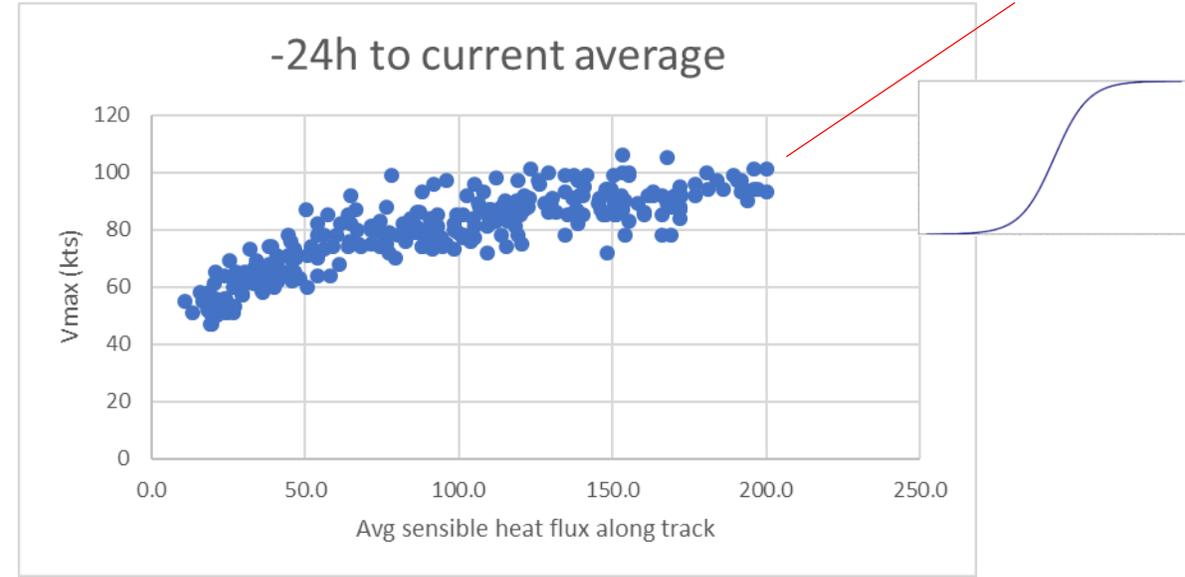
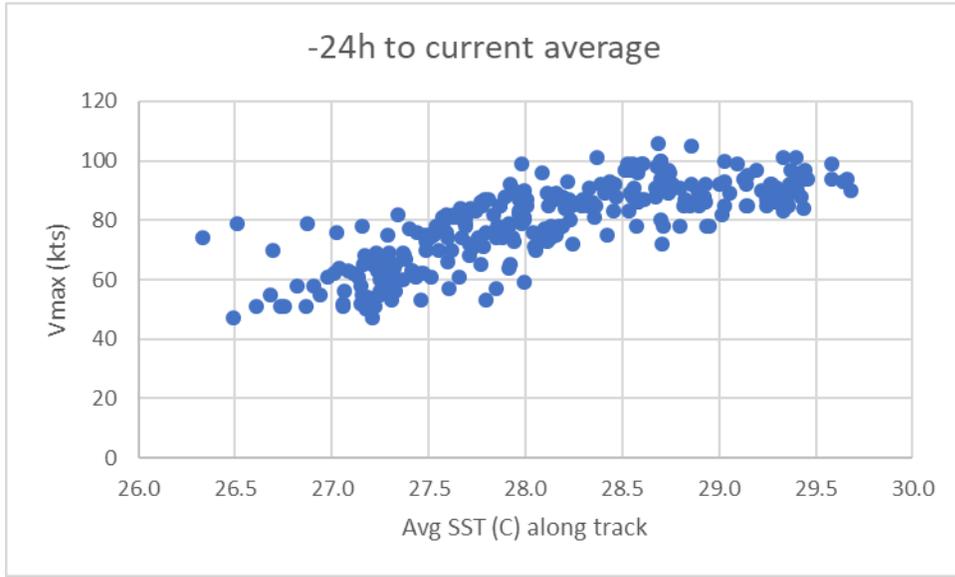


# Relationship analysis

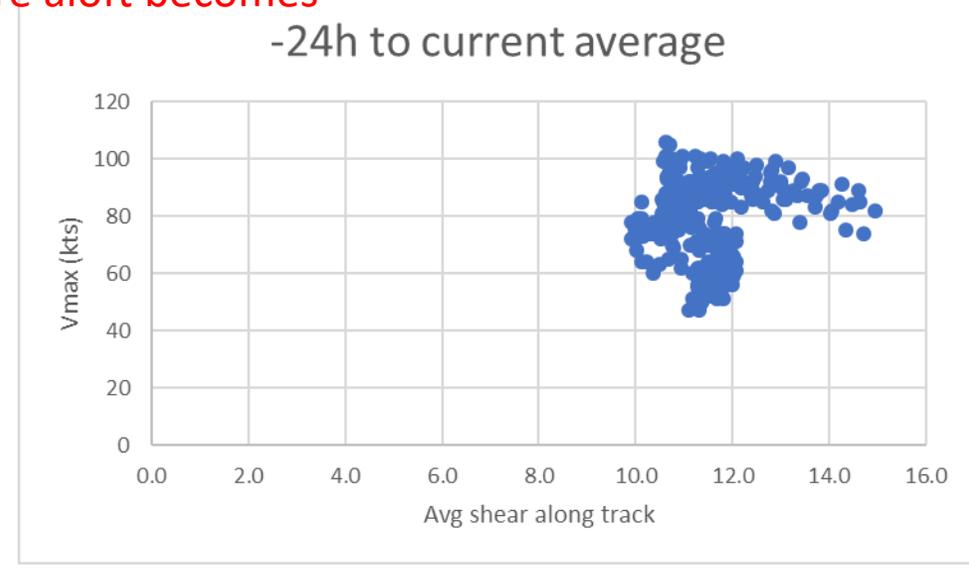
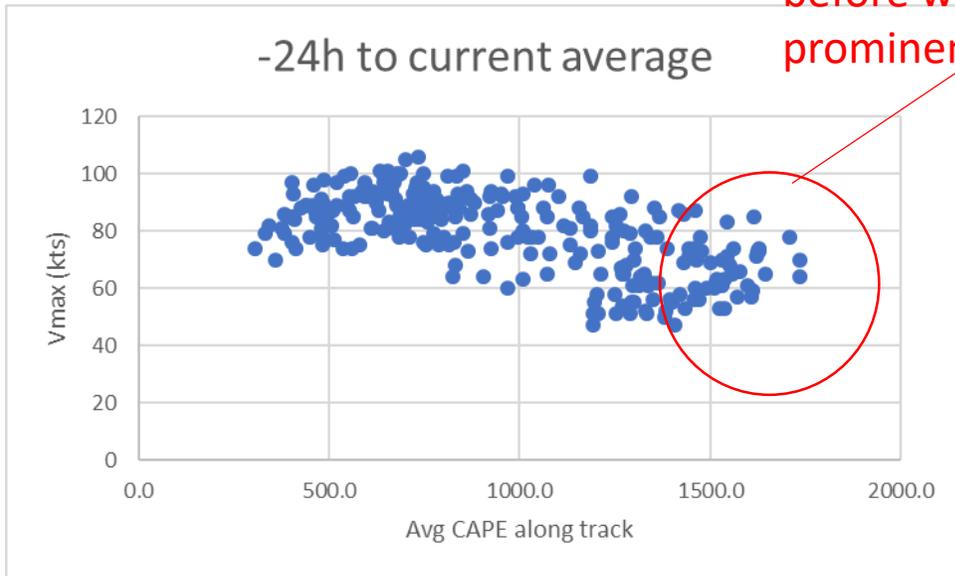




Sigmoidal relationship



Associated with SST=27.5-28.5C, and before warm core aloft becomes prominent





## Fitted equation for CAPE, constrained by heat flux sigmoidal relationship

Integrate CAPE linear regression equation

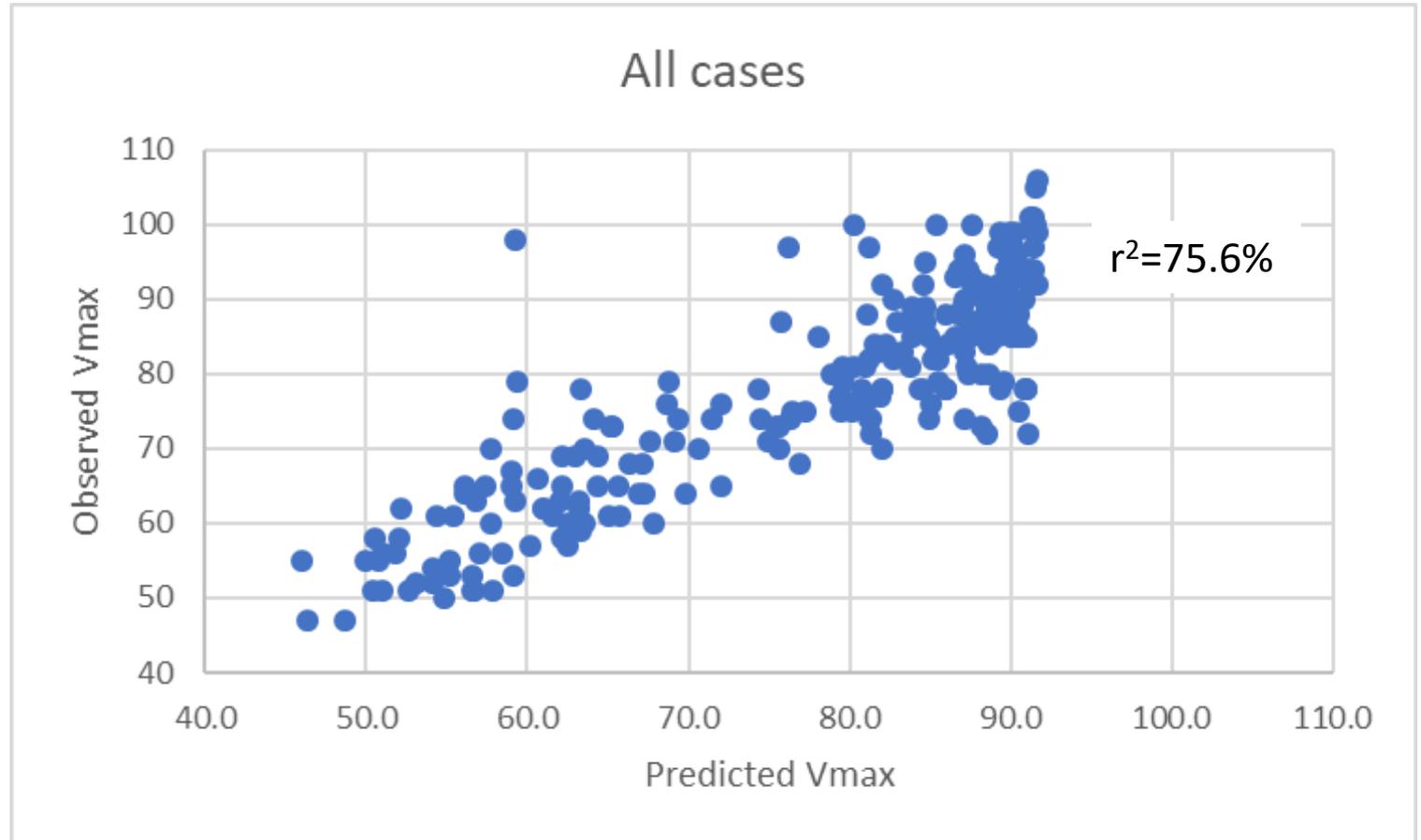
$$V_{max}(24) = V_{max}(0) - 10.2 + 0.02\overline{CAPE}$$

If  $\overline{CAPE} < 75$  J/kg, or  $V_{max}(H) < V_{max}(24)$ , then

$$V_{max}(H) = \frac{92.34}{1 + e^{-(H-19.77)/39.17}}$$

In a simplistic nutshell, possibly explains HWRF sensitivity to SST in a favorable PW, low-to-moderate shear environment.

CAPE boost from increasing SST (in this study 27.5-28.5C, but probably environment- and storm-specific) results in faster intensity change. As CAPE is depleted, steady-state occurs depending on critical flux thresholds.



Maximum Potential Intensity studies?

Assume general sigmoidal relationship

$$V_{max}(H) = \frac{\varepsilon MPI}{1+e^{-(H-A)/B}} = \frac{\varepsilon MPI}{1+\frac{e^{A/B}}{e^{H/B}}}$$

where  $H$  is sensible heat flux,  $\varepsilon$  is an environmental inhibitor ( $\varepsilon = 1$  is MPI conditions), and  $A$ ,  $B$  are empirically-derived constants.

Assume bulk equation for  $H = \rho C_E c_p V \Delta T$ , where  $\Delta T = T_{air} - SST$   
 $\rho = 1.18 \text{ kg/m}^3$ ,  $C_E = 0.00118$ ,  $c_p = 1004 \text{ J/kg-K}$ . One obtains an implicit equation:

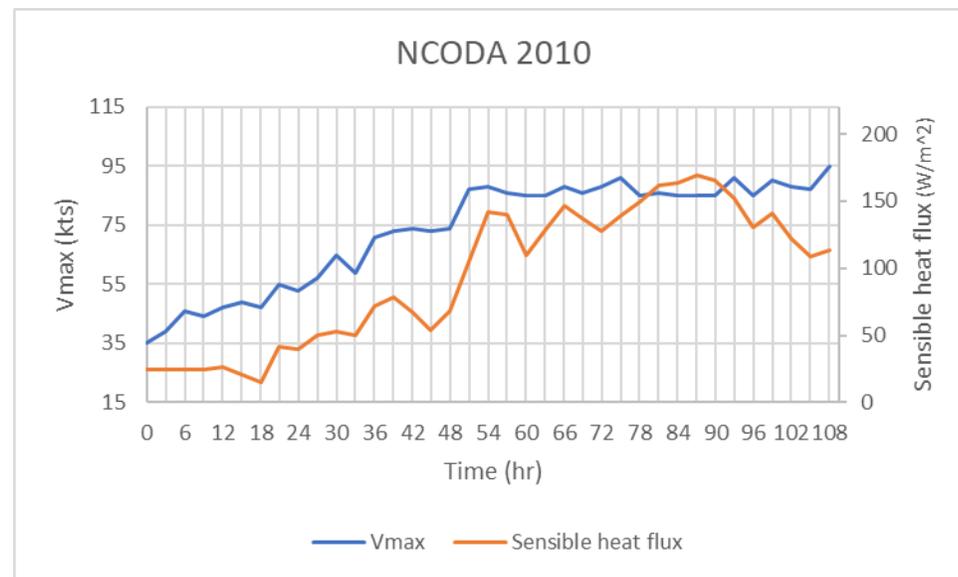
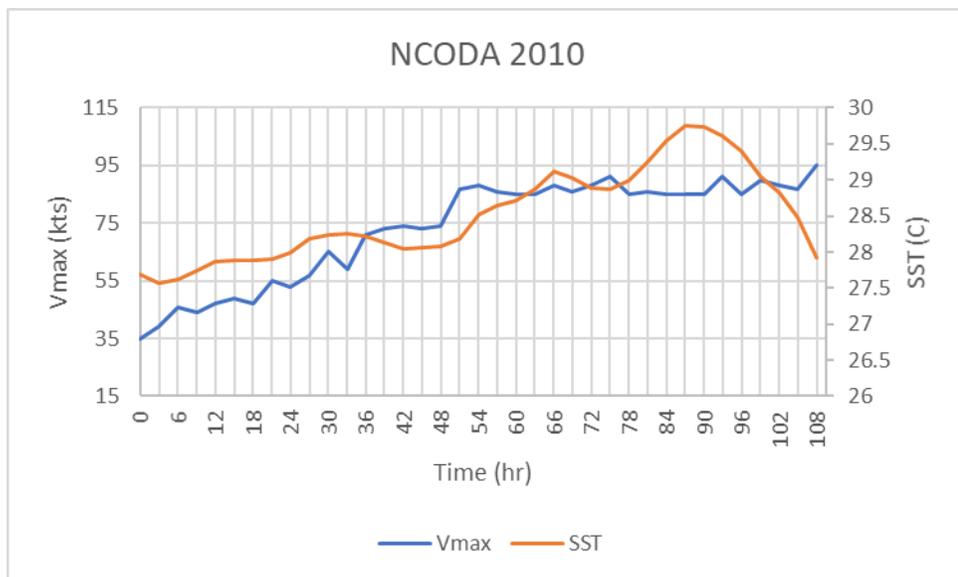
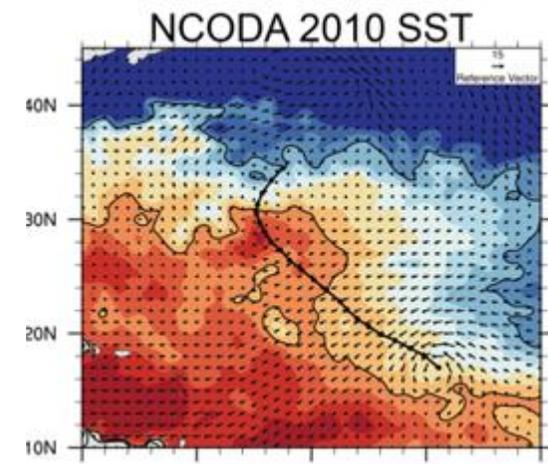
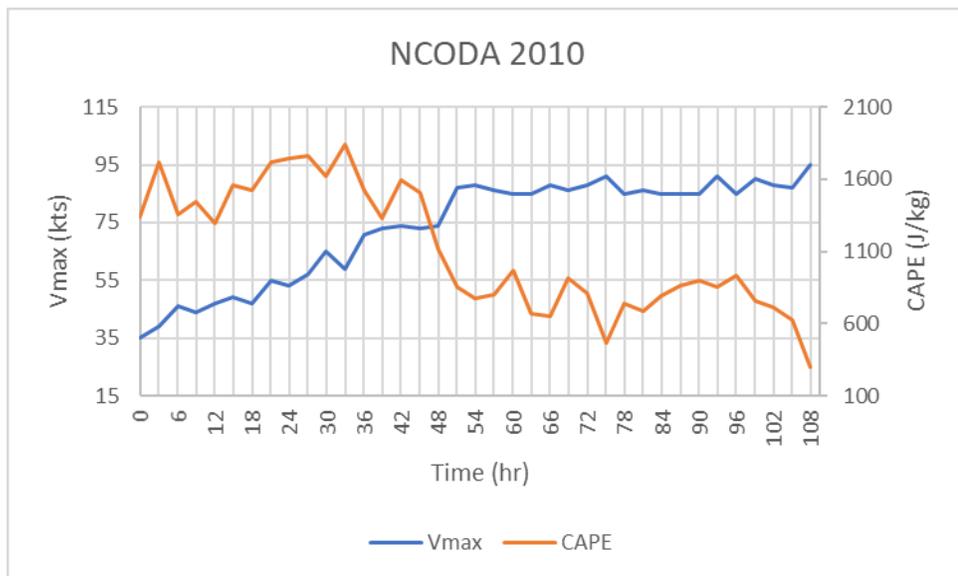
$$V_{max}(H) = \frac{\varepsilon MPI}{1+\exp\left(-\frac{1.4V_{max}\Delta T+A}{B}\right)} = \frac{\varepsilon MPI}{1+\frac{\exp\left(\frac{A}{B}\right)}{\exp\left(\frac{1.4V_{max}\Delta T}{B}\right)}}$$

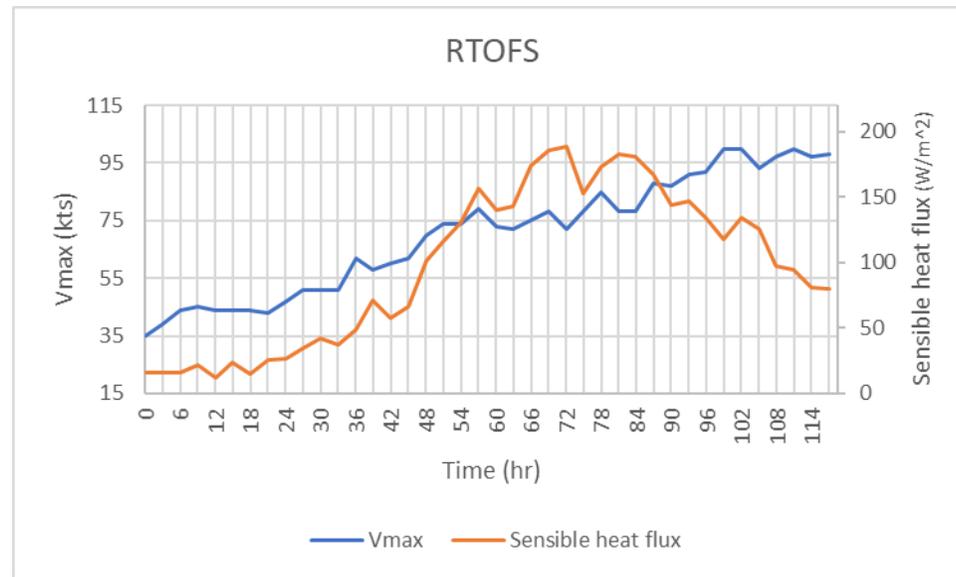
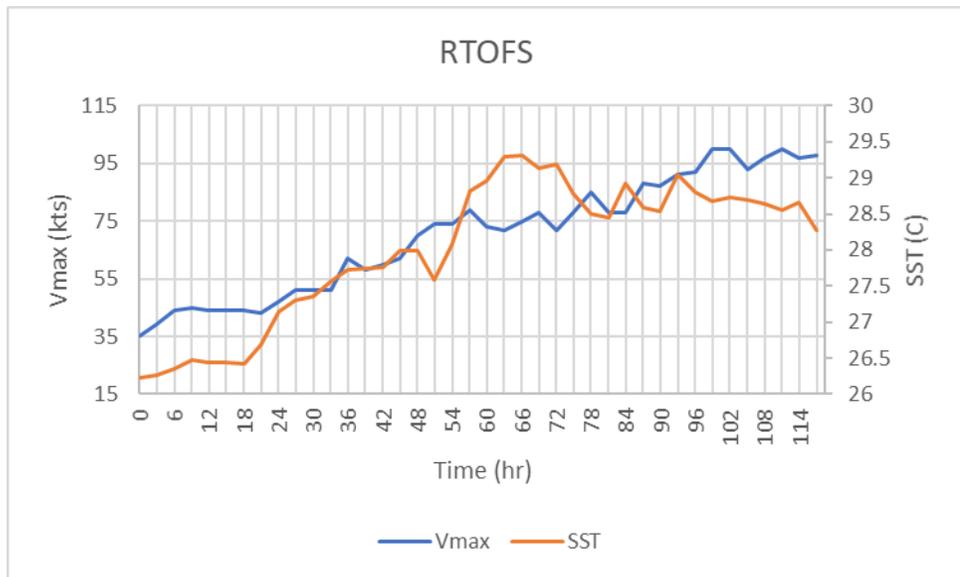
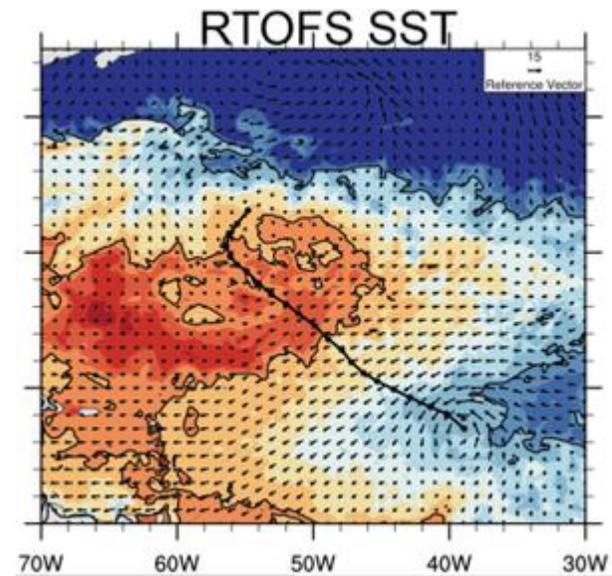
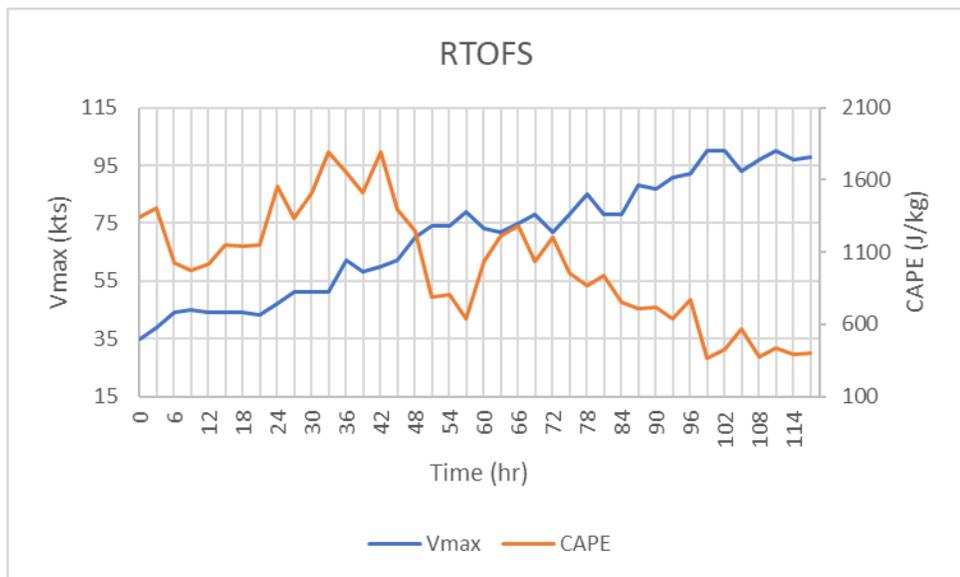
A large HWRF SST-sensitivity database could elucidate steady-state and MPI functionality with surface fluxes through empirical application.

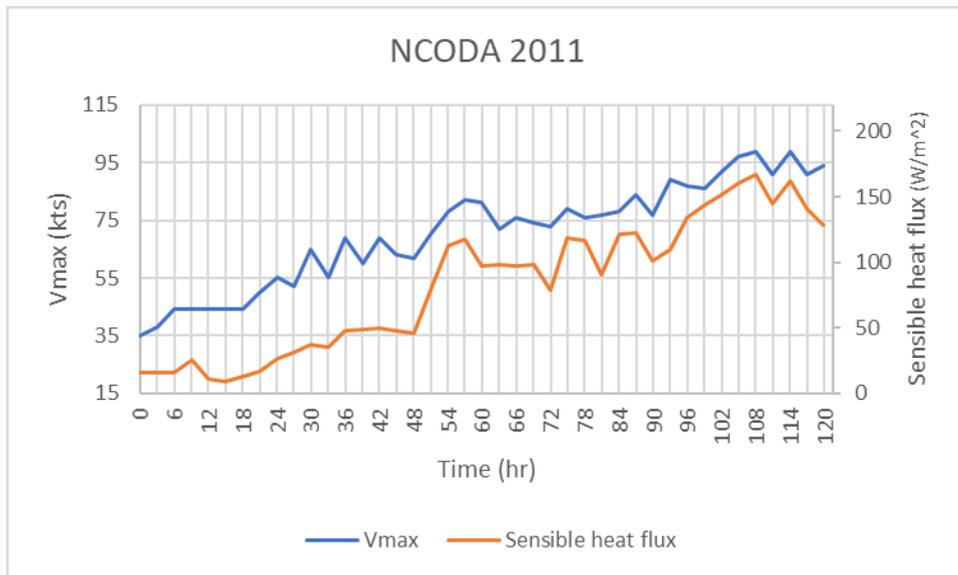
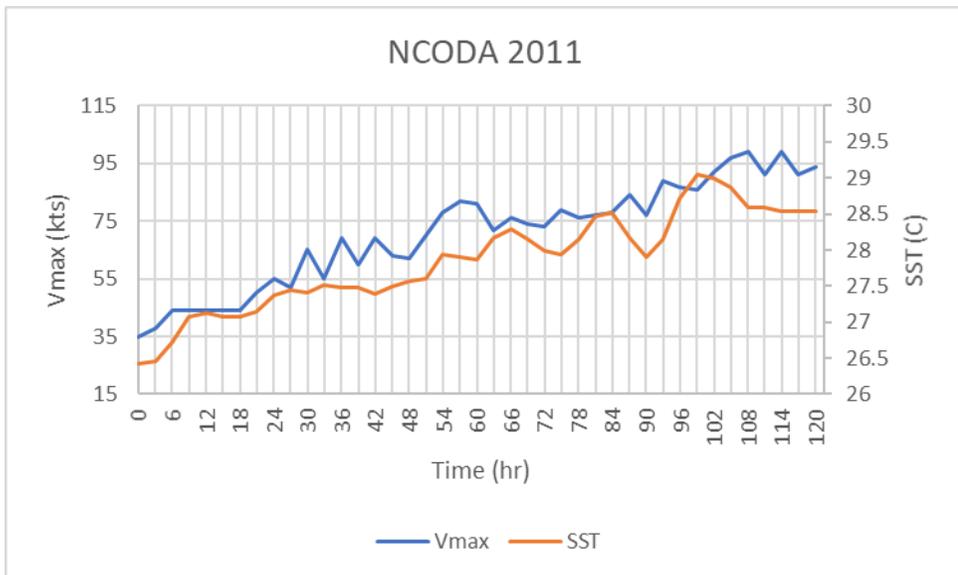
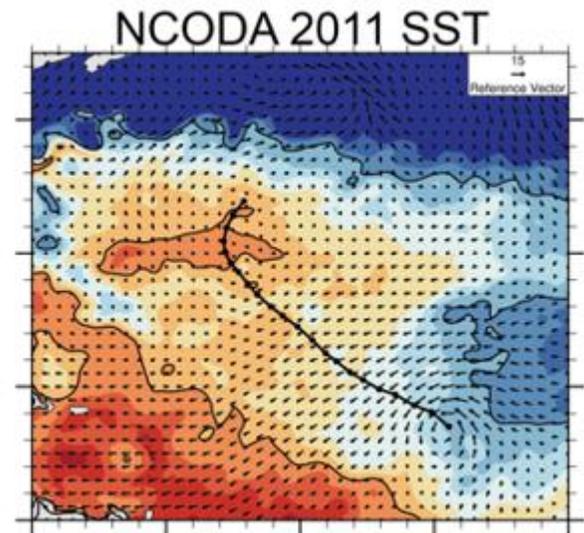
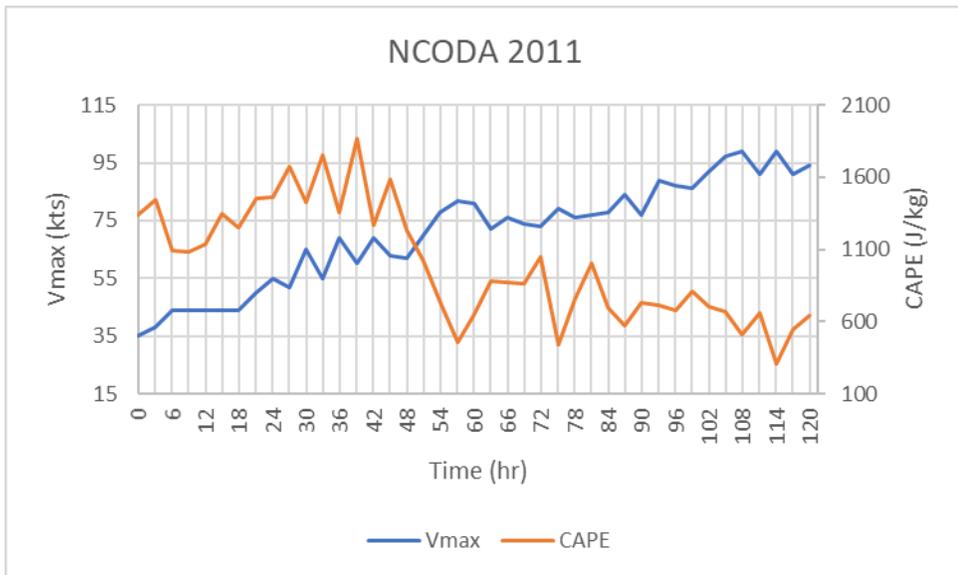
# Recommendations

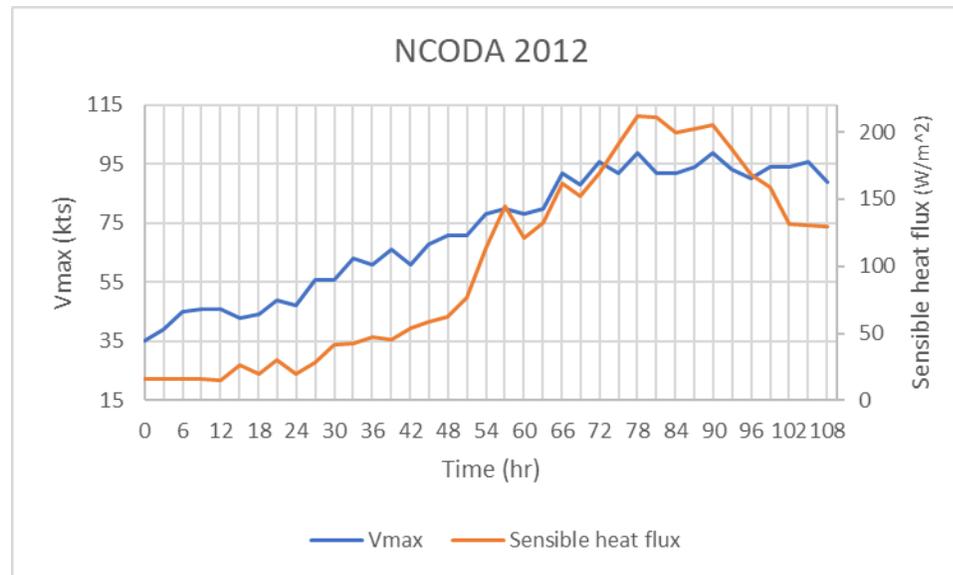
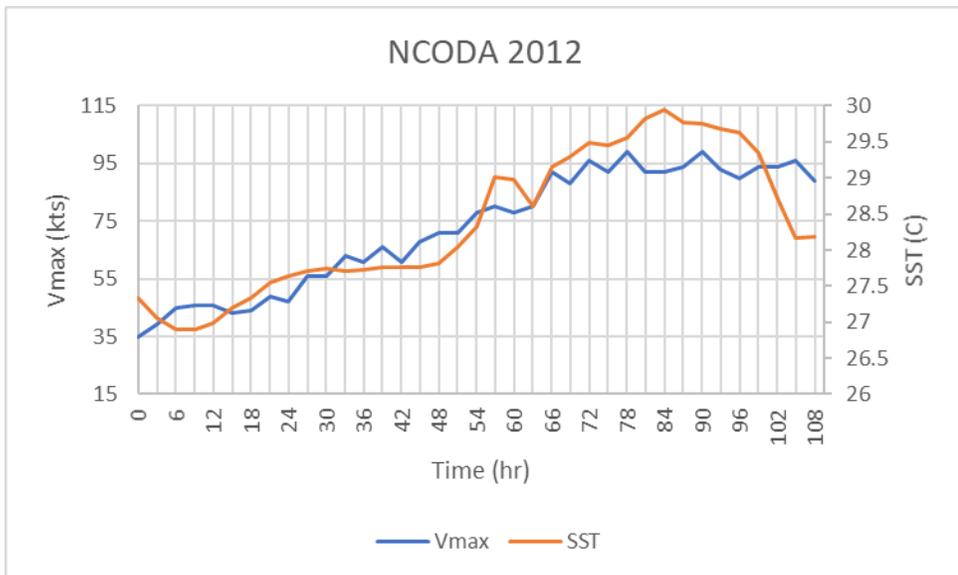
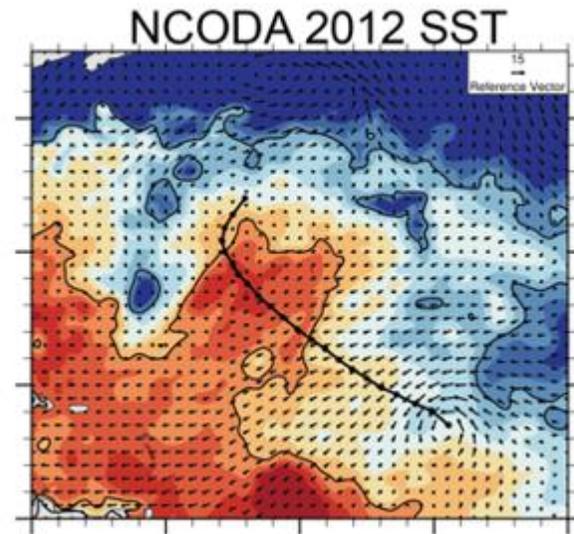
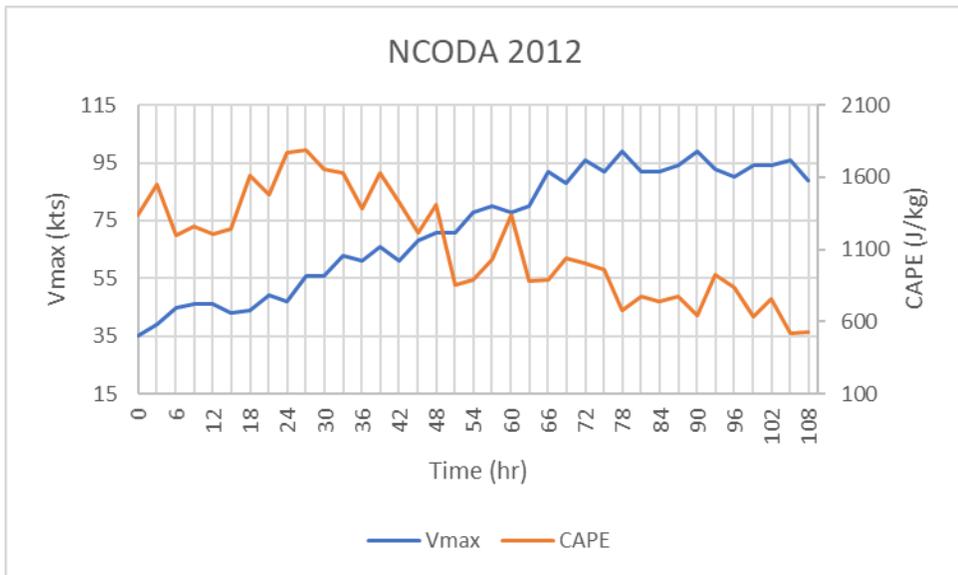
- 1) HWRF intensity skill has exceeded SHIPS and LGEM two years in a row. Its plausible a large database of HWRF simulations could now glean an understanding of rapid intensification situations.
- 2) Suggests using model CAPE may be useful in statistical schemes such as SHIPS, LGEM, and Rapid Intensification Index.
- 3) Studies will also be useful for understanding Maximum Potential Intensity theory.
- 4) Study supports the Ocean Model Impact Tiger Team (OMITT) emphasis on surface forcing processes to improve intensity forecasts. Events such as Harvey's unexpectedly rapid intensification, possibly due to an under-analyzed warm pool, illustrate the need for better coupling processes.

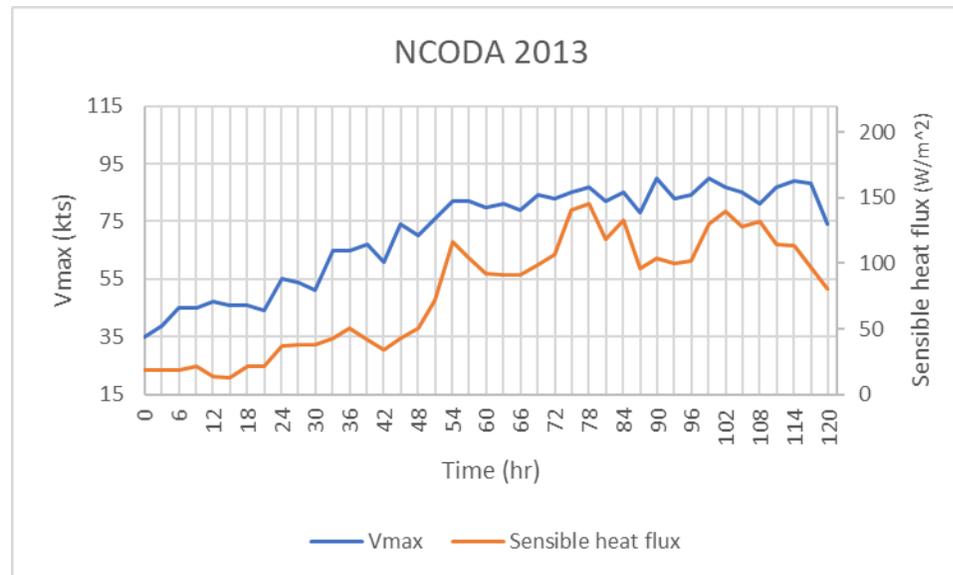
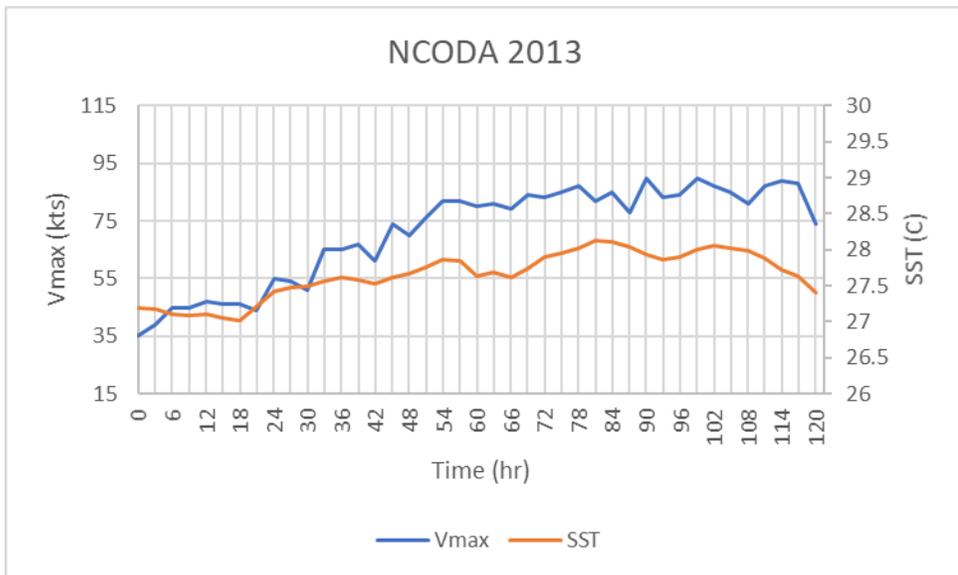
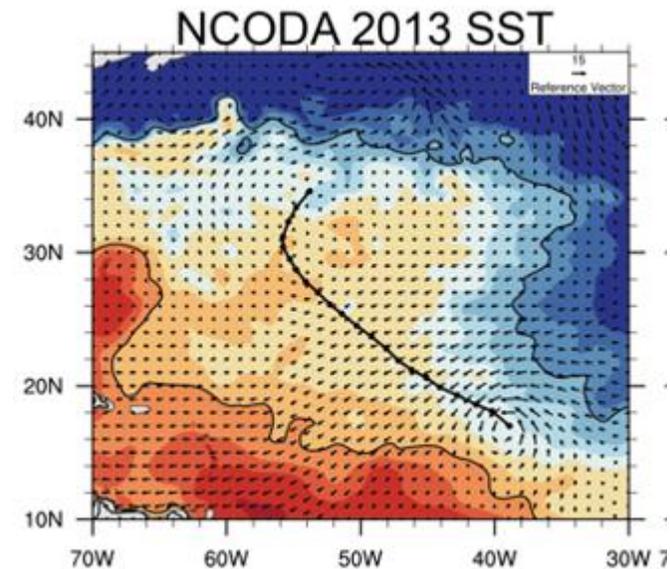
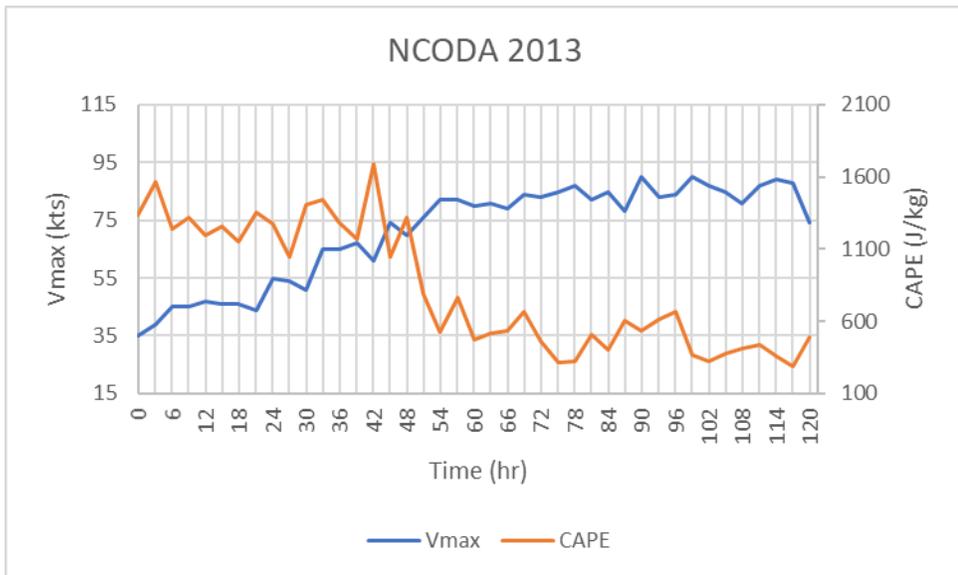
Extra slides



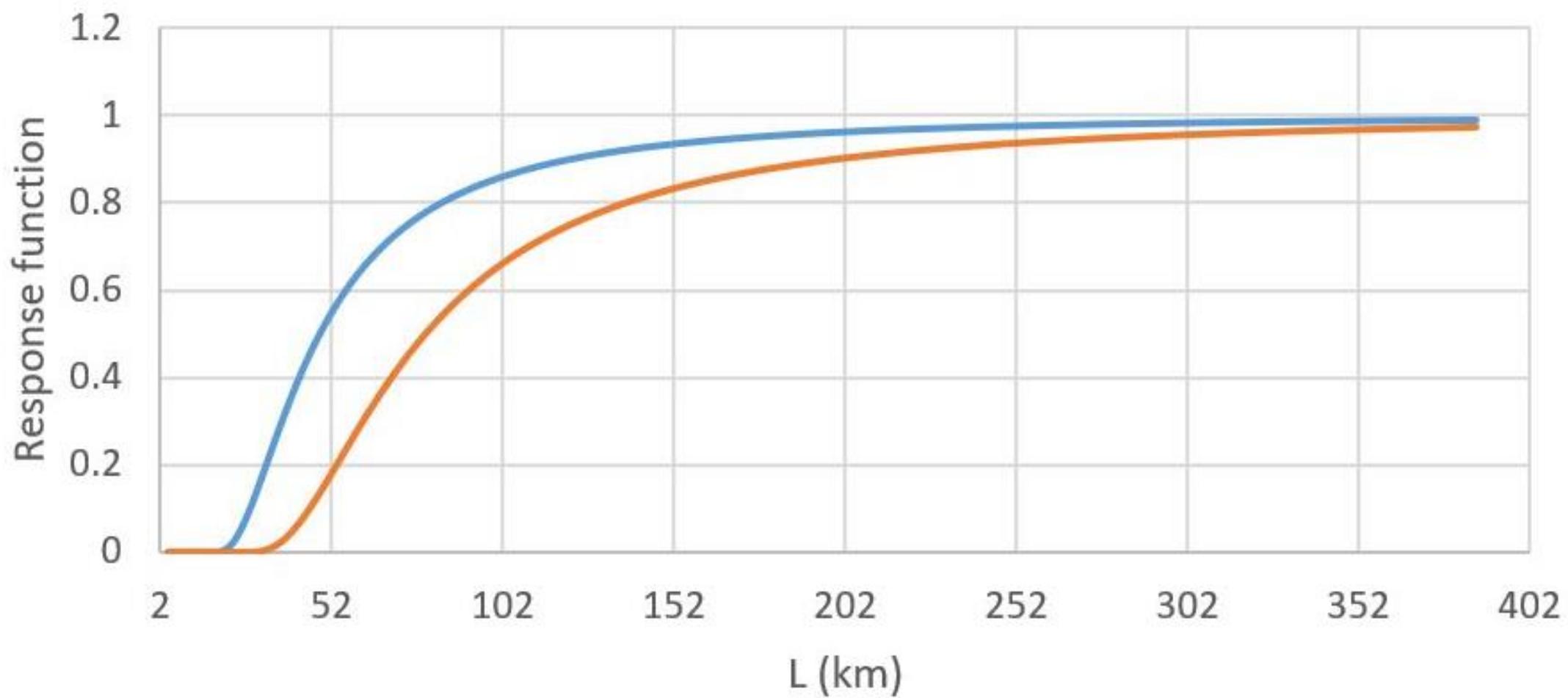








## Attenuation per pass



## Spatial filter analysis

This type of equation tries to remove the shortest waves but leaves the longer ones relatively unaffected, and has analogies to diffusion

$$\textcircled{1} \quad \phi_i^{\tilde{t}+1} = \phi_i^{\tilde{t}} + K \left( \phi_{i+1}^{\tilde{t}} + \phi_{i-1}^{\tilde{t}} - 2\phi_i^{\tilde{t}} \right)$$

The properties of this filter can be explored by representing waves in terms of complex variables

$$\textcircled{2} \quad \phi(x,t) = \hat{\phi}(k,\omega) e^{i(kx+\omega t)}$$

where  $\hat{\phi}$ ,  $k$ , and  $\omega$  can be complex and when discretized:

$$x = n\Delta x, \quad n = \pm 1, \pm 2, \dots$$

$$t = \tilde{t}\Delta t, \quad \tilde{t} = 0, 1, 2, \dots \text{ either as a timestep or \# passes}$$

$\omega = \text{frequency}$

$k = \text{wavenumber}$

Denote:

$$\phi_i^{\tilde{t}+1} = \phi_i^{\tilde{t}} \phi_0^1; \quad \phi_{i+1}^{\tilde{t}} = \phi_i^{\tilde{t}} \phi_1^0; \quad \phi_{i-1}^{\tilde{t}} = \phi_i^{\tilde{t}} \phi_{-1}^0$$

Substitute (2) into (1), and cancel common  $\hat{\phi}_i$  and  $\hat{\phi}$  factors

$$(3) \quad e^{i\omega\Delta t} = 1 + K \left( e^{ik\Delta x} + e^{-ik\Delta x} - 2 \right)$$

Decompose into real and imaginary parts  $\omega = \omega_r + i\omega_i$

$= 2 \cos(k\Delta x)$

$$e^{i\omega\Delta t} = e^{i(\omega_r + i\omega_i)\Delta t} = e^{i\omega_r\Delta t} e^{-\omega_i\Delta t}$$

because of  $i^2 = -1$

(4) Denote  $\lambda = \pm e^{-\omega_i\Delta t} \equiv$  amplitude change of the solution per time step or for each pass.

Use Euler's formula

$$(5) \quad e^{i\omega_r\Delta t} = \cos(\omega_r\Delta t) + i \sin(\omega_r\Delta t)$$

Substitute (4) and (5) into (3)

$$\lambda \left( \cos(\omega_r\Delta t) + i \sin(\omega_r\Delta t) \right) = 1 + 2K \left( \cos(k\Delta x) - 1 \right)$$

Solve for real and imaginary parts

$$\lambda \cos(\omega_r\Delta t) = 1 + 2K \left( \cos(k\Delta x) - 1 \right)$$

$$\lambda \sin(\omega_r\Delta t) = 0$$

Second equation shows that  $\omega_r = 0$ , and there is no propagation. Also, then  $\cos(\omega_r \Delta x) = 1$

Therefore

$$\lambda = 1 + 2K(\cos(k\Delta x) - 1)$$

Since  $k = \frac{2\pi}{L}$  where  $L \equiv \text{wavelength}$ , and  $L = n\Delta x$

⑥  $\lambda = 1 + 2K\left(\cos\left(\frac{2\pi}{n}\right) - 1\right)$

For a value of  $K = \frac{1}{4}$

$n$	2	3	4	6	8	10	15	20	32
$\lambda$	0	0.25	0.5	0.75	0.85	0.9	0.96	0.975	0.99

Hence the 20x wave is removed, the 30x-60x is extremely attenuated, and longer wavelengths are retained.

But suppose  $K = 1.707107$ , corresponding to  $K = 0.5 / (1 - \cos(\frac{2\pi}{m}))$  where  $m = 8$

$n$	2	3	4	6	8	10	15	20	32
$\lambda$	-5.8	-4.1	-2.4	-0.71	0	0.35	0.71	0.83	0.93

$n=2$  to  $n=5$  are unstable,  $n=6$  to  $n=7$  are negatively attenuated,  $n=8$  is removed,  $n=9$  to  $n=16$  are extremely attenuated, and longer wavelengths are retained

Hence, if  $m=n$ , the max wavelength is removed  
if  $m > n$ , the max wavelength is either negatively attenuated or unstable.

Generally, if  $n > m+3$ , attenuation is minimal

Generally, if  $m > n+3$ , unstable (but depends on  $n$ , can be  $n+2$  or  $n+4$ )

Kurihara sequentially tries to remove  $n=2$  to  $n=9$ , but because new zax waves are generated by round-off error for  $m=3$  to  $m=9$ , to keep scheme stable and to remove the zax waves, he periodically applies a  $m=2$

Note also that as model grid spacing decreases, the effectiveness of the original Kurihara scheme will decrease. This is why Fitzpatrick and Lay added new weights, shown in red.

$m = 2, 3, 4, 2, 5, 6, 7, 2, 8, 9, 2, 10, 11, 2, 12, 13, 2, 3, 2$